Letter to the Editor

# On the relationship between the fundamental matrices for different definitions of the state vector 

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## 1. Introduction

Recently, the need arose, in the context of a research work, to know the relationship between the fundamental matrices of a linear mechanical system with $n$ degrees-of-freedom (d.o.f.), the state vector of which is usually defined in two different forms. As is known, some authors define it as the $(2 n \times 1)$ vector of the $n$ generalized co-ordinates and the $n$ generalized velocities [1-3], whereas some others define it in the reverse order [4,5]. The desired relationship was not found in the literature. Although it is acknowledged that the contribution of this study does not solve a very complex problem, it is nevertheless thought that the simple result established in the present letter can be helpful for those working in this area.

## 2. Theory

As is known, the free vibrations of a discrete linear mechanical system with $n$ d.o.f. is governed in the physical space by the following matrix differential equation of order two:

$$
\begin{equation*}
\mathbf{M} \ddot{\boldsymbol{q}}+\mathbf{D} \dot{\mathbf{q}}+\mathbf{K} \mathbf{q}=0, \tag{1}
\end{equation*}
$$

where M, D and $\mathbf{K}$ denote the $(n \times n)$ mass, damping and stiffness matrices, respectively. $\mathbf{q}(t)$ represents the $(n \times 1)$ vector of the generalized co-ordinates to describe the position of the mechanical system .This differential equation can equivalently be written in the so-called statespace form

$$
\begin{equation*}
\dot{\mathbf{x}}=\mathbf{A x} \tag{2}
\end{equation*}
$$

[^0]where the $(2 n \times 1)$ state vector $\mathbf{x}(t)$ and the $2 n \times 2 n$ system matrix $\mathbf{A}$ are defined as
\[

$$
\begin{gather*}
\mathbf{x}=\left[\mathbf{q}^{\mathrm{T}} \cdots \dot{\mathbf{q}}^{\mathrm{T}}\right]^{\mathrm{T}}  \tag{3}\\
\mathbf{A}=\left[\begin{array}{ccc}
\mathbf{0} & \vdots & \mathbf{I} \\
\cdots & \cdots & \cdots \\
-\mathbf{M}^{-1} \mathbf{K} & \vdots & -\mathbf{M}^{-1} \mathbf{D}
\end{array}\right] \tag{4}
\end{gather*}
$$
\]

I being the ( $n \times n$ ) unit matrix.
The general solution of the differential Eq. (2) can be written as [2]

$$
\begin{equation*}
\mathbf{x}(t)=\boldsymbol{\Phi}(t) \mathbf{x}_{0} \tag{5}
\end{equation*}
$$

where $\mathbf{x}_{0}$ denotes the initial state vector and the $(2 n \times 2 n)$ fundamental matrix is defined as

$$
\begin{equation*}
\boldsymbol{\Phi}(t)=\mathbf{X e}^{\Lambda t} \mathbf{X}^{-1} \tag{6}
\end{equation*}
$$

where the modal matrix $\mathbf{X}$ and $\mathrm{e}^{1 t}$ are defined as

$$
\begin{equation*}
\mathbf{X}=\left[\tilde{\mathbf{x}}_{1}, \ldots, \tilde{\mathbf{x}}_{2 n}\right], \quad \mathrm{e}^{\Lambda t}=\boldsymbol{\operatorname { d i a g }}\left(\mathrm{e}^{\lambda_{j} t}\right) \quad(j=1, \ldots, 2 n) \tag{7}
\end{equation*}
$$

Here $\lambda_{j}$ and $\tilde{\mathbf{x}}_{j}$ denote the $j$ th eigenpair corresponding to the eigenvalue problem regarding (2).
Now, let it be assumed that the state vector is defined as

$$
\begin{equation*}
\mathbf{x}^{\prime}=\left[\dot{\mathbf{q}}^{\mathrm{T}}: \mathbf{q}^{\mathrm{T}}\right]^{\mathrm{T}} \tag{8}
\end{equation*}
$$

such that the state-space equation reads now

$$
\begin{equation*}
\dot{\mathbf{x}}^{\prime}=\mathbf{A}^{\prime} \mathbf{x}^{\prime} \tag{9}
\end{equation*}
$$

The counterparts of Eqs. (4)-(6) are

$$
\begin{gather*}
\mathbf{A}^{\prime}=\left[\begin{array}{ccc}
-\mathbf{M}^{-1} \mathbf{D} & \vdots & -\mathbf{M}^{-1} \mathbf{K} \\
\cdots & \cdots & \cdots \\
\mathbf{I} & \vdots & \mathbf{0}
\end{array}\right]  \tag{10}\\
\mathbf{x}^{\prime}(t)=\boldsymbol{\Phi}^{\prime}(t) \mathbf{x}_{0}^{\prime}  \tag{11}\\
\mathbf{\Phi}^{\prime}(t)=\mathbf{X}^{\prime} \mathrm{e}^{\Lambda t} \mathbf{X}^{\prime-1} \tag{12}
\end{gather*}
$$

where the modal matrix $\mathbf{X}^{\prime}$ consists of the new $2 n$ eigenvectors $\tilde{\mathbf{x}}_{j}^{\prime}$ of the eigenvalue problem, corresponding to (9).

The following ( $2 n \times 2 n$ ) matrix be introduced

$$
\mathbf{Q}=\left[\begin{array}{ccc}
\mathbf{0} & \vdots & \mathbf{I}  \tag{13}\\
\cdots & \cdots & \cdots \\
\mathbf{I} & \vdots & \mathbf{0}
\end{array}\right]
$$

It can easily be shown that $\mathbf{Q}$ is orthogonal and is equal to its own inverse. It can further be shown that the modal matrices of both representations are interrelated by

$$
\begin{equation*}
\mathbf{X}^{\prime}=\mathbf{Q X} \tag{14}
\end{equation*}
$$

If this is substituted into Eq. (12),

$$
\begin{equation*}
\boldsymbol{\Phi}^{\prime}(t)=\mathbf{Q} \mathbf{X} \mathrm{e}^{\Lambda t} \mathbf{X}^{-1} \mathbf{Q} \tag{15}
\end{equation*}
$$

is obtained. Considering Eq. (6), the last formula reduces to

$$
\begin{equation*}
\boldsymbol{\Phi}^{\prime}(t)=\mathbf{Q} \boldsymbol{\Phi}(t) \mathbf{Q} \tag{16}
\end{equation*}
$$

which represents the desired result.

## References

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